

# Modern Introductory Physics, Part II, Exam 3

Friday, Mar. 29, 2024 — Covering Moore Chapters 10-12.

YOU MAY NEED SOME SEPARATE PAPER TO WORK THE PROBLEMS. AT THE END, PLEASE STAPLE YOUR WORK TO THE EXAM AND TURN IT ALL IN.

## 1. Building a Table of Gaussian Integrals

Let's repeat and extend some of the things you did on Problem Set 10. First we need to rebuild our table of integrals.

(a) Start with  $\int_{-\infty}^{\infty} e^{-cx^2} dx = \sqrt{\pi/c}$ . Treating  $c$  as the variable, take  $\frac{d}{dc}$  of each side of this equation and simplify to get a formula for  $\int_{-\infty}^{\infty} x^2 e^{-cx^2} dx$ .

(b) Take another derivative with respect to  $c$  and simplify to get a formula for  $\int_{-\infty}^{\infty} x^4 e^{-cx^2} dx$ .

(c) Summarize the results so you can use them in the next problem:

$$\int_{-\infty}^{\infty} e^{-cx^2} dx = \sqrt{\pi/c}$$

$$\int_{-\infty}^{\infty} x^2 e^{-cx^2} dx =$$

$$\int_{-\infty}^{\infty} x^4 e^{-cx^2} dx =$$

(d) Summarize the results again, but substituting  $c = 1/b^2$ ,

$$\int_{-\infty}^{\infty} e^{-x^2/b^2} dx =$$

$$\int_{-\infty}^{\infty} x^2 e^{-x^2/b^2} dx =$$

$$\int_{-\infty}^{\infty} x^4 e^{-x^2/b^2} dx =$$

## 2. Normalization and Expectation Values

It can be important in determining whether quantum mechanics predictions are right (it is too much to be more specific here), to know things like  $\overline{x^4} = \int_{-\infty}^{\infty} \psi^{*}(x) x^4 \psi(x) dx = \int_{-\infty}^{\infty} x^4 |\psi(x)|^2 dx$ , or as an even more important example,  $\overline{p^4} = \int_{-\infty}^{\infty} \psi^{*}(x) \left(\frac{\hbar}{i} \frac{d}{dx}\right)^4 \psi(x) dx$ .

(a) For the ground state of the harmonic oscillator,  $\psi(x) = c_0 e^{-x^2/2b^2}$ . Use your results from 1(d) to normalize  $\psi(x)$ , e.g., to determine  $c_0$  by demanding  $1 = \int_{-\infty}^{\infty} |\psi(x)|^2 dx$ .

(b) Use your results from 1(d) and 2(a) to determine

$$\overline{x^2} = \int_{-\infty}^{\infty} x^2 |\psi(x)|^2 dx$$

(c) Use your results from 1(d) and 2(a) to determine

$$\overline{x^4} = \int_{-\infty}^{\infty} x^4 |\psi(x)|^2 dx$$

(ASIDE: For those of you who have had or will have statistics, you need  $\bar{x} = \int_{-\infty}^{\infty} x P(x) dx$ ,

$\overline{x^2} = \int_{-\infty}^{\infty} x^2 P(x) dx$ ,  $\overline{x^3} = \int_{-\infty}^{\infty} x^3 P(x) dx$ , and  $\overline{x^4} = \int_{-\infty}^{\infty} x^4 P(x) dx$  to calculate the mean, variance, skew, and kurtosis of a distribution.)

### 3. Reflection at a Potential Jump

We are going to learn about reflection at jump in the potential.

(a) Imagine a de Broglie wave coming in from the left (e.g., from negative  $x$ ) and hitting a barrier at  $x = 0$ .

$$\psi_L(x, t) = e^{i k_L x - i \omega t}$$

Using the time-dependent Schrödinger equation,  $i\hbar \frac{\partial \psi_L(x, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_L(x, t)}{\partial x^2}$ , what is the formula for  $k_L$  in terms of  $\hbar$ ,  $\omega$ , and  $m$ ?

(b) We haven't normalized  $\psi_L(x, t)$  and we can't. It represents an infinite stream of particles coming at the barrier from the left. There might be a reflected stream too. We can add this to  $\psi_L(x, t)$  as follows:

$$\psi_L(x, t) = e^{i k_L x - i \omega t} + b e^{-i k_L x - i \omega t}$$

What does  $\psi_L(x, t)$  simplify to at  $x = 0$ ?

(c) Take  $\frac{\partial \psi_L(x, t)}{\partial x}$ .

(d) What does  $\frac{\partial \psi_L(x, t)}{\partial x}$  simplify to at  $x = 0$ ?

## 4. Transmission at a Potential Jump

This is actually a continuation of Problem 3. If there is a barrier, and it isn't too high, we expect there to be a transmitted wave going to the right.

(a) We can model that as

$$\psi_R(x, t) = c e^{i k_R x - i \omega t}$$

We'll have the barrier be  $V_0$  high for  $x > 0$ , and we'll assume  $\hbar\omega > V_0$ .

(ASIDE: We assume  $\hbar\omega > V_0$  because that is what "isn't too high" amounts to. We could alternatively assume  $\hbar\omega < V_0$ , in which case the region  $x > 0$  would be classically forbidden. But that is a different problem.)

Using the time-dependent Schrödinger equation,  $i\hbar \frac{\partial \psi_R(x, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_R(x, t)}{\partial x^2} + V_0 \psi_R(x, t)$ , what is the formula for  $k_R$  in terms of  $\hbar$ ,  $\omega$ ,  $m$  and  $V_0$ ?

(b) What does  $\psi_R(x, t)$  simplify to at  $x = 0$ ?

(c) Take  $\frac{\partial \psi_R(x, t)}{\partial x}$ .

(d) What does  $\frac{\partial \psi_R(x, t)}{\partial x}$  simplify to at  $x = 0$ ?

(e) Using your answers for 3(b) and 4(b) what does continuity demand?

(f) Using your answers for 3(d) and 4(d) what does the no-kinks condition demand?

HINT/CROSS-CHECK: If you did the algebra and simplification right, for parts (e) and (f) you will have two equations for the two unknowns  $b$  and  $c$ . The equations will not involve  $x$  or  $t$  or  $\omega$ . They just involve  $k_L$  and  $k_R$  (which in turn depend on all the other constants). If we wanted to continue and learn more about transmission and reflection, next we would solve these equations for  $b$  and  $c$ .

## 5. Harmonic Oscillators in Thermal Equilibrium

You got a taste of fermions in thermal equilibrium in the section on semiconductors. In this problem, we will work on the simpler situation of a lot of harmonic oscillators in thermal equilibrium.

To make life easier, instead of writing  $\frac{1}{k_B T}$  a lot, we will define  $\beta \equiv \frac{1}{k_B T}$ .

The fundamental principle of statistical mechanics is that for a system in thermal equilibrium, the probability,  $P_n$ , that a system has energy  $E_n$  is proportional to  $e^{-\beta E_n}$ . We don't know the proportionality constant yet. I will just write it as  $\frac{1}{Z}$ , so  $P_n = \frac{1}{Z} e^{-\beta E_n}$ .

(a) For the harmonic oscillator  $E_n = \hbar\omega(n + \frac{1}{2})$ . What is  $\frac{P_1}{P_0}$ ?

(b) What is  $\frac{P_{n+1}}{P_n}$ ?

(c) To determine  $Z$ , it has to be that the sum of all the  $P_n$  adds up to 1. so  $1 = \sum_{n=0}^{\infty} \frac{1}{Z} e^{-\beta E_n}$  or

$Z = \sum_{n=0}^{\infty} e^{-\beta E_n}$ . Put in  $E_n = \hbar\omega(n + \frac{1}{2})$ . Factor out the overall  $e^{-\frac{1}{2}\beta\hbar\omega}$ . You are close to knowing  $Z$  for the harmonic oscillator.

(d) Use  $e^{-\beta\hbar\omega n} = (e^{-\beta\hbar\omega})^n$  and  $\sum_{n=0}^{\infty} x^n = \frac{1}{1-x}$  to finalize your answer for  $Z$ .

(e) Put your answer for  $Z$  into  $P_n = \frac{1}{Z} e^{-\beta E_n}$  and simplify. When you are done simplifying, you will know  $P_n$  in terms of the constants  $\beta$ ,  $\hbar$ ,  $\omega$ ,  $n$ , and of course  $e$ .

Name \_\_\_\_\_

1. /5

2. /5

3. /5

4. /5

5. /5

GRAND TOTAL

/25